

Rigidity and flexibility of entropies of boundary maps associated to Fuchsian groups

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Set-up and history

- $\Gamma :=$ cocompact torsion free Fuchsian group (surface group),
 $\Gamma < \text{PSU}(1, 1)$ acting properly discontinuously by
orientation-preserving isometries on the unit disk \mathbb{D} .

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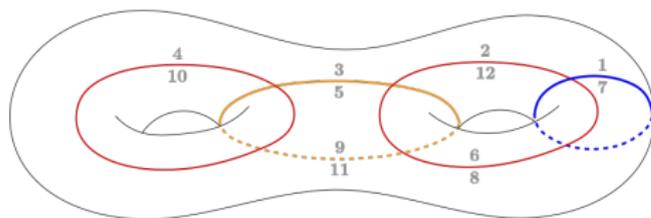
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 - **1991: Adler-Flatto** studied **Bowen-Series maps** for surface groups using fundamental $8g - 4$ -gons.
 - **2017: S.K.-Ugarcovici, 2019: Abrams-S.K.** studied families of boundary maps $f_{\tilde{A}}$ for surface groups related to fundamental $8g - 4$ -gons and proved that 2-dim their natural extensions $F_{\tilde{A}}$ have global **attractors with finite rectangular structure**.

$(8g - 4)$ -gonal fundamental domain \mathcal{F}

Any compact orientable hyperbolic surface $S = \Gamma \backslash \mathbb{D}$ of genus $g \geq 2$ admits an $(8g - 4)$ -gonal fundamental domain $\mathcal{F} \subset \mathbb{D}$ obtained by cutting it with $2g$ closed geodesics that intersect in pairs (g of them go around the “holes” and another g go around the “waists” of S).

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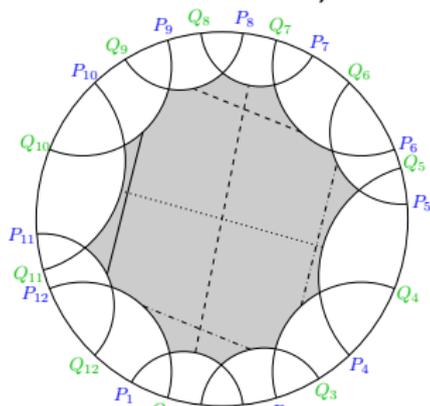
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$g = 2$

Identification of sides by the involution σ :

$$\sigma(k) := \begin{cases} 4g - k \pmod{8g - 4} & \text{if } k \text{ is odd} \\ 2 - k \pmod{8g - 4} & \text{if } k \text{ is even} \end{cases}$$



$g = 2$

Dehn, Fenchel, Nielsen, Koebe, Adler-Flatto [AF91].

[AF91] R. Adler, L. Flatto. Geodesic flows, interval maps, and symbolic dynamics, *Bull. Amer. Math. Soc.* **25** (1991), No. 2, 229–334.

Marked canonical polygon

Thus obtained $8g - 4$ -gon \mathcal{F} does not need to be regular, but satisfies the following properties:

- 1 The identified sides k and $\sigma(k)$ have equal length.
 - 2 The angles at glued vertices k and $\sigma(k) + 1$ add up to π
- and is called **marked canonical polygon**.

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$\Gamma = \langle T_k \rangle, \{T_k : \text{side } k \rightarrow \text{side } \sigma(k)\}$ is a Fuchsian group of the first kind with signature $(g; -; 0)$, i.e. a surface group.

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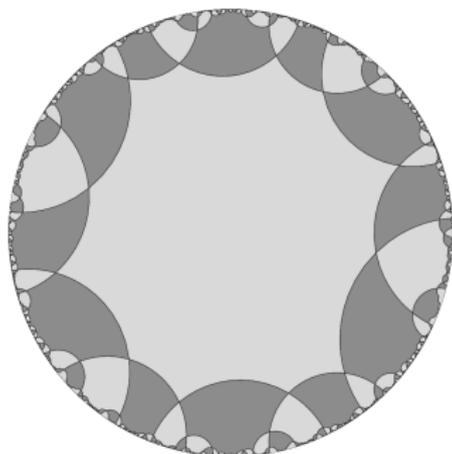
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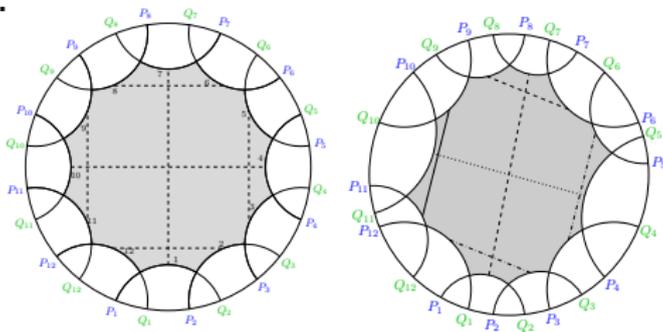
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Teichmüller space $\mathcal{T}(g)$ via fundamental polygons

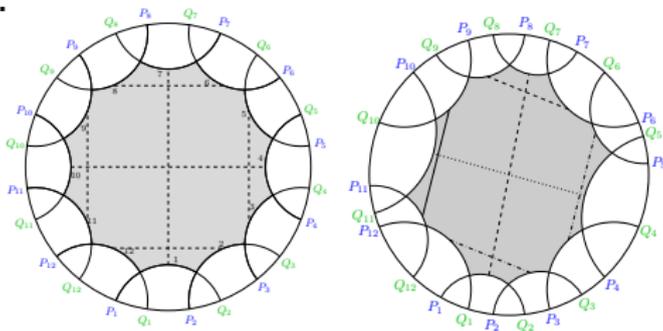
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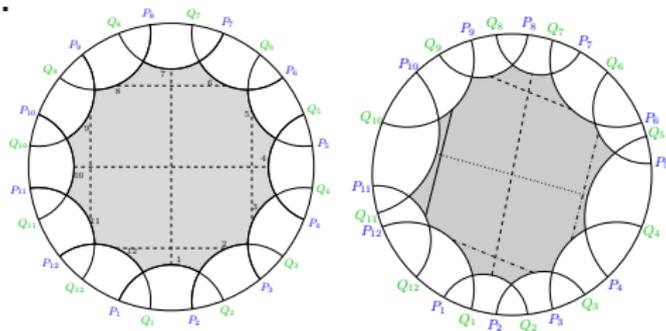
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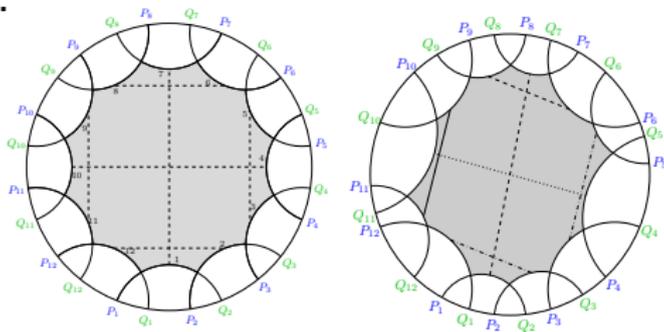
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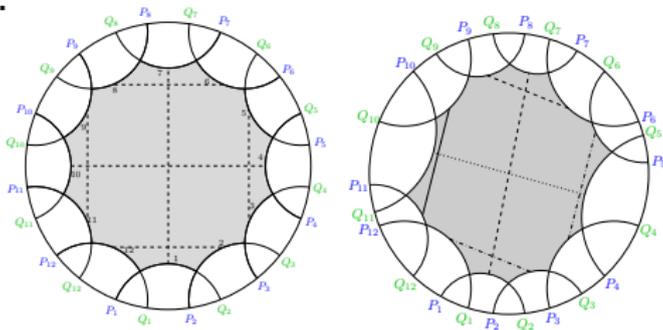
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- This construction gives a less common representation of the Teichmüller space $\mathcal{T}(g)$ for $g \geq 2$: it is **the space of marked canonical $(8g - 4)$ -gons in \mathbb{D} up to an isometry of \mathbb{D} .**

A family of boundary maps

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- For any marked canonical fundamental $8g - 4$ -gon $\mathcal{F} \in \mathcal{T}(g)$, let $\bar{P} = \{P_1, \dots, P_{8g-4}\}$, $\bar{Q} = \{Q_1, \dots, Q_{8g-4}\}$, P_k and Q_{k+1} are the endpoints of the geodesic containing side k ;

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- For any **multi-parameter** $\bar{A} = \{A_1, \dots, A_{8g-4}\}$, $A_k \in [P_k, Q_k]$ we define the **boundary map** $f_{\bar{A}} : \mathbb{S} \rightarrow \mathbb{S}$

$$f_{\bar{A}}(x) = T_k(x) \quad \text{if } x \in [A_k, A_{k+1}),$$

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$f_{\bar{P}}$ was referred to in [AF91] as the **Bowen–Series boundary map**, (although Bowen–Series construction [BS79] used $4g$ -sided polygons).

[BS79] R. Bowen, C. Series. Markov maps associated with Fuchsian groups, *Inst. Hautes Études Sci. Publ. Math.* No. 50 (1979), 153–170.

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Deformations

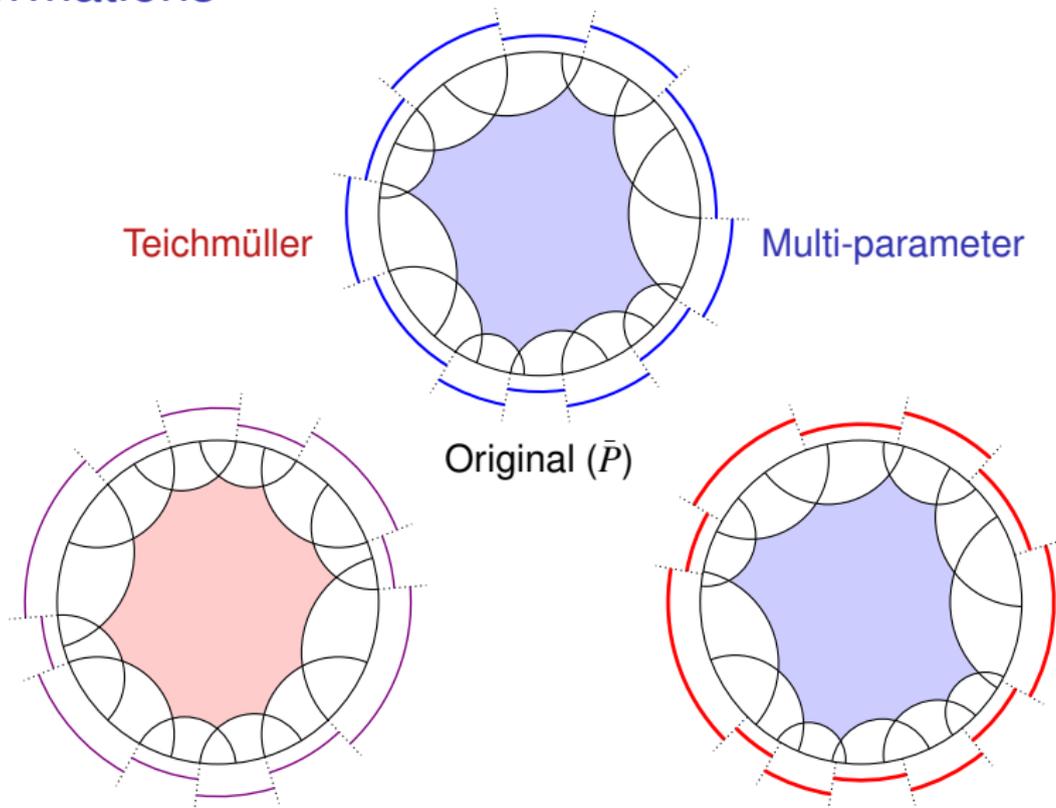
Teichmüller

Multi-parameter

Original (\bar{P})

Change polygon

Change interval where
each T_k is used



Multi-parameters

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e.g. \bar{P} and \bar{Q} are extremal.
- [KU17]: If $A_k \in (P_k, Q_k)$ \bar{A} satisfies the **cycle property**, i.e. \exists positive integers m_k, n_k such that $f^{m_k}(T_k A_k) = f^{n_k}(T_{k-1} A_k)$

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 - Each $f_{\bar{A}}$ is a **piecewise monotone piecewise continuous** map of the circle.
 - If \bar{A} is **extremal or satisfies the short cycle property**, then $f_{\bar{A}}$ admits a **unique smooth invariant ergodic measure $\mu_{\bar{A}}$ related to the Liouville measure for the geodesic flow**.

Rigidity and Flexibility theorems

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Let $S = \Gamma \backslash \mathbb{D}$ be a surface of genus $g \geq 2$ and let $\mathcal{F} \in \mathcal{T}(g)$.

Theorem 1. [AKU-2024] Rigidity of topological entropy.

$\forall \bar{A} = \{A_1, \dots, A_{8g-4}\}$ with $A_k \in [P_k, Q_k]$, the map $f_{\bar{A}} : \mathbb{S} \rightarrow \mathbb{S}$ has the same topological entropy $h_{\text{top}}(f_{\bar{A}}) = \log(4g - 3 + \sqrt{(4g - 3)^2 - 1})$.

Note: $f_{\bar{A}}$ are generally not topologically conjugate since, according to [KU17], the combinatorial structure of orbits associated to A_k could differ.

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Theorem 2. [AKU-2022] Flexibility of measure-theoretic entropy

(in the spirit of flexibility program suggested by Anatole Katok). Let \bar{A} be extremal or satisfy the short cycle property (an open set of partitions). Then

- 1 For each $\mathcal{F} \in \mathcal{T}(g)$, $h_{\mu_{\bar{A}}}(f_{\bar{A}}) = \frac{\pi^2(4g-4)}{\text{Perimeter}(\mathcal{F})} = \pi \cdot \frac{\text{Area}(\mathcal{F})}{\text{Perimeter}(\mathcal{F})}$.
- 2 The maximum value of $h_{\mu_{\bar{A}}}(f_{\bar{A}})$, $H(g) = \frac{\pi^2(4g-4)}{(8g-4) \operatorname{arccosh}(1+2 \cos \frac{\pi}{4g-2})}$ is achieved on the surface with \mathcal{F}_{reg} .
- 3 For each $t \in (0, H(g)] \exists \mathcal{F} \in \mathcal{T}(g)$ such that $h_{\mu_{\bar{A}}}(f_{\bar{A}}) = t$.

[AKU-2024] A.Abrams, S.Katok, I.Ugarcovici, In *A Vision for Dynamics - The Legacy of Anatole Katok*, 19 - 47, Cambridge University Press, 2024

[AKU-2022] A.Abrams, S.Katok, I.Ugarcovici, *Ergod. Th. & Dynam. Sys.* **42** (2022), 389 - 401

Topological entropy of one-dimensional maps: history

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- Alsedà-Misiurewicz (2015): piecewise continuous non-Markov

Steps of the proof of Theorem 1

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① Topological entropy for Markov cases: extremal \bar{A} . Since

$f_{\bar{A}}(P_k), f_{\bar{A}}(Q_k) \in \bar{P} \cup \bar{Q}$, partition of \mathbb{S} into intervals I_1, \dots, I_{16g-8} :

$I_{2k-1} := [P_k, Q_k], \quad I_{2k} := [Q_k, P_{k+1}], \quad k = 1, \dots, 8g - 4,$

for any extremal \bar{A} , $f_{\bar{A}}$ is Markov with $M_{\bar{A}} = (m_{i,j})$ over the alphabet

$\{1, \dots, 16g - 8\}$, $m_{i,j} := \begin{cases} 1 & \text{if } f_{\bar{A}}(I_i) \supset I_j \\ 0 & \text{otherwise} \end{cases}$.

Below are $M_{\bar{P}}$ (left) and $M_{\bar{Q}}$ (right) for $g = 2$.

$$\begin{pmatrix} 00000000000000000000000000000000 \\ 1111111111100000000011111111 \\ 011000000000000000000000000000 \\ 0001111111111111111111100000 \\ 0000000000001100000000000000 \\ 11111100000000111111111111 \\ 000000000000000000000000110 \\ 111111111111111111000000001 \\ 00000001100000000000000000 \\ 11000000011111111111111111 \\ 00000000000000000001100000 \\ 11111111111100000000011111 \\ 00011000000000000000000000 \\ 0000011111111111111111100 \\ 000000000000011000000000 \\ 111111110000000111111111 \\ 1000000000000000000000001 \\ 011111111111111111000000 \\ 000000000110000000000000 \\ 111100000001111111111111 \\ 00000000000000000000011000 \\ 111111111111100000000111 \\ 000001100000000000000000 \\ 000000011111111111111111 \end{pmatrix} \begin{pmatrix} 11000000000000000000000000 \\ 11111111111000000001111111 \\ 00000000001100000000000000 \\ 0001111111111111111110000 \\ 00000000000000000000001100 \\ 111111000000001111111111 \\ 00000011000000000000000000 \\ 1111111111111111000000001 \\ 0000000000000000011000000 \\ 110000000111111111111111 \\ 00110000000000000000000000 \\ 11111111111100000000011111 \\ 00000000000011000000000000 \\ 000001111111111111111100 \\ 000000000000000000000011 \\ 111111110000000011111111 \\ 00000000110000000000000000 \\ 011111111111111111000000 \\ 0000000000000000000110000 \\ 111100000001111111111111 \\ 00001100000000000000000000 \\ 111111111111100000000111 \\ 000000000000000110000000 \\ 000000011111111111111111 \end{pmatrix}$$

- The number of \bar{A} -admissible words of length n grows \approx as λ^n , where $\lambda = 4g - 3 + \sqrt{(4g - 3)^2 - 1}$, \Rightarrow the maximal eigenvalue of every $M_{\bar{A}}$ is λ .

Steps of the proof of Theorem 1

① Topological entropy for Markov cases: extremal \bar{A} . Since

$f_{\bar{A}}(P_k), f_{\bar{A}}(Q_k) \in \bar{P} \cup \bar{Q}$, partition of \mathbb{S} into intervals I_1, \dots, I_{16g-8} :

$$I_{2k-1} := [P_k, Q_k], \quad I_{2k} := [Q_k, P_{k+1}], \quad k = 1, \dots, 8g - 4,$$

for any extremal \bar{A} , $f_{\bar{A}}$ is Markov with $M_{\bar{A}} = (m_{i,j})$ over the alphabet

$$\{1, \dots, 16g - 8\}, m_{i,j} := \begin{cases} 1 & \text{if } f_{\bar{A}}(I_i) \supset I_j \\ 0 & \text{otherwise} \end{cases}.$$

Below are $M_{\bar{P}}$ (left) and $M_{\bar{Q}}$ (right) for $g = 2$.

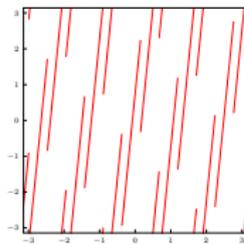
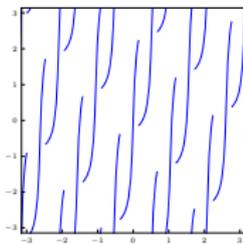
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- The number of \bar{A} -admissible words of length n grows \approx as λ^n , where $\lambda = 4g - 3 + \sqrt{(4g - 3)^2 - 1}$, \Rightarrow the maximal eigenvalue of every $M_{\bar{A}}$ is λ .
- Thus if \bar{A} is extremal, $h_{\text{top}}(f_{\bar{A}}) = \log \lambda$, so $h_{\text{top}}(f_{\bar{P}}) = h_{\text{top}}(f_{\bar{Q}}) = \log \lambda$.

② Conjugacy to a constant slope map: $\ell_{\bar{A}} = \psi_{\bar{A}} \circ f_{\bar{A}} \circ \psi_{\bar{A}}^{-1}$.

$f_{\bar{P}}$

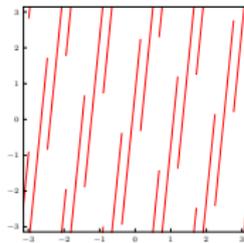
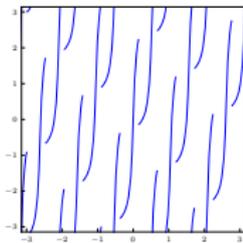
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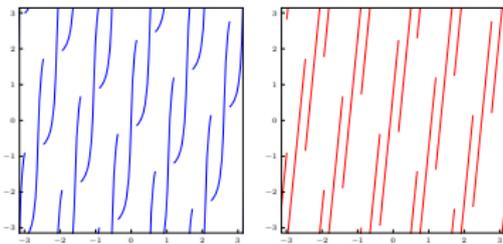


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This is the most technically difficult part of the argument (explicit formulae for $\psi_{\bar{P}}$ and $\psi_{\bar{Q}}$ via second Parry measure, relation between Markov partitions and cylinder intervals for $f_{\bar{P}}$ and $f_{\bar{Q}}$: re-coding).

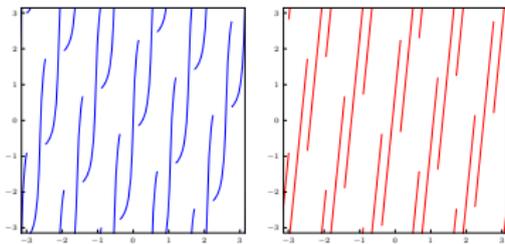


For $g = 2$, $I_{\bar{P}}^{(1,16)}$ as a union of 17 \bar{Q} -cylinders (left) and $I_{\bar{P}}^{(1,17)}$ as a union of 2 \bar{Q} -cylinders (right)

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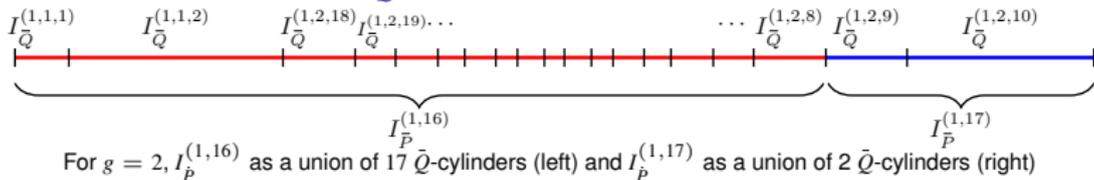
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③ Completion of the proof: the conjugacy $\psi_{\bar{P}}$ works for all $\bar{A} = \{A_1, \dots, A_{8g-4}\}$ with $A_k \in [P_k, Q_k]$. $\psi_{\bar{P}} \circ f_{\bar{A}} \circ \psi_{\bar{P}}^{-1}$ is piecewise linear with the same slope λ , hence by [M-S 1980] $h_{\text{top}}(f_{\bar{A}}) = \log \lambda$.

Steps of the proof of Theorem 2

① The smooth invariant measure.

In coordinates (u, w, s) , where s is the arclength along geodesic from u to w , the Liouville volume on $T^1\mathbb{D}$ can be written as $d\omega = 2dm$, where $dm = d\nu ds$, $d\nu = \frac{|du| |dw|}{|u - w|^2}$ is the smooth measure on the space of oriented geodesics \mathbb{D} modeled as $\mathbb{S} \times \mathbb{S} \setminus \Delta$, called **geodesic current**.

☛ We defined a 2-dim map $F_{\bar{A}} : \mathbb{S} \times \mathbb{S} \setminus \Delta \rightarrow \mathbb{S} \times \mathbb{S} \setminus \Delta$ by

$$F_{\bar{A}}(u, w) = (T_k u, T_k w) \quad \text{if } w \in [A_k, A_{k+1}).$$

[KU17]: $F_{\bar{A}}$ has a global attractor $\Omega_{\bar{A}}$ with finite rectangular structure.

☛ The restriction of $F_{\bar{A}}$ to $\Omega_{\bar{A}}$ is the **natural extension map** of $f_{\bar{A}}$.

☛ $F_{\bar{A}} : \Omega_{\bar{A}} \rightarrow \Omega_{\bar{A}}$ preserves the smooth probability measure $d\nu_{\bar{A}} = \frac{d\nu}{\int_{\Omega_{\bar{A}}} d\nu}$.

☛ The geodesic flow on S φ^t is a special flow over a cross-section parametrized by $\Omega_{\bar{A}}$, s.t. the first return map to this cross-section is

$F_{\bar{A}} : \Omega_{\bar{A}} \rightarrow \Omega_{\bar{A}}$ [AK19].

☛ $f_{\bar{A}}$ is a factor of $F_{\bar{A}}$ (projecting on the second coordinate), so one can obtain its smooth invariant probability measure $\mu_{\bar{A}}$ on \mathbb{S} as a projection.

② Formula for the entropy.

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Using that the entropy of φ^t with respect to the normalized Liouville measure is equal to 1, Abramov's formula and Ambrose-Kakutani Theorem we obtained in [AK19] the formula

$$h_{\mu_{\tilde{A}}}(f_{\tilde{A}}) = h_{\mu_{\tilde{A}}}(F_{\tilde{A}}) = \frac{\pi^2(4g - 4)}{\int_{\Omega_{\tilde{A}}} d\nu} = \pi \cdot \frac{\text{Area}(\mathcal{F})}{\int_{\Omega_{\tilde{A}}} d\nu}.$$

(The last equality uses Gauss–Bonnet formula: $\text{Area}(\mathcal{F}) = 2\pi(2g - 2)$.)

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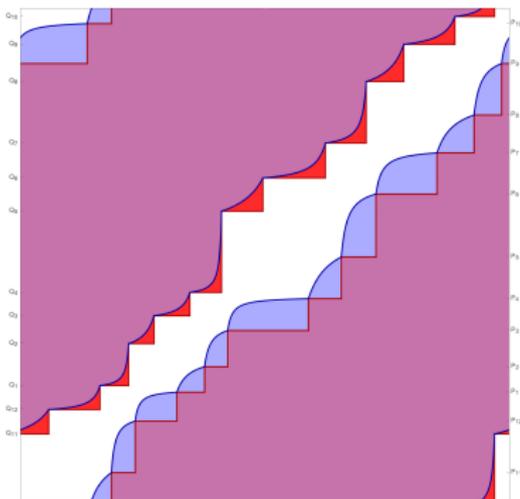
• **A new ingredient:** $\int_{\Omega_{\tilde{A}}} d\nu = \text{Perimeter}(\mathcal{F})$. We use another map, also introduced by Adler–Flatto in [AF91], called in [AK19] **geometric map**. It is defined on

$$\Omega_G := \{ (u, w) : uw \text{ intersects } \mathcal{F} \} \subset \mathbb{S} \times \mathbb{S} \setminus \Delta \text{ by}$$

$$F_G(u, w) = (T_k u, T_k w) \quad \text{if } uw \text{ exits } \mathcal{F} \text{ through side } k.$$

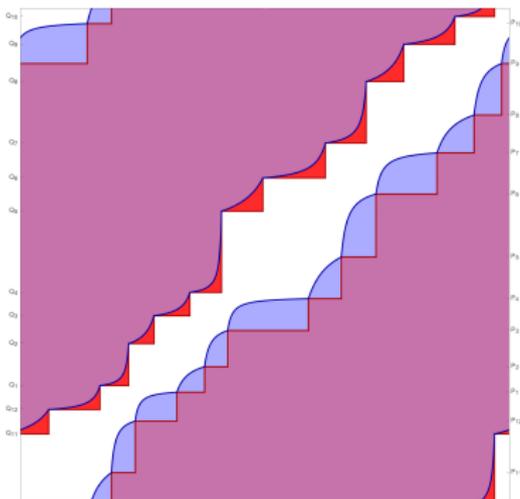
[AK19] A. Abrams, S. Katok. Adler and Flatto revisited: cross-sections for geodesic flow on compact surfaces of constant negative curvature. *Studia Mathematica* **246** (2019), 167–202.

Theorem [AK19, [A20]: \exists is a bijection map $\Phi : \Omega_G \rightarrow \Omega_{\bar{A}}$,
 piecewise-Möbius which maps **bulges** into **corners**.



Arithmetic set Ω_p in red-and-purple, geometric set Ω_G in blue-and-purple (for an irregular polygon with $g = 2$)

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Corollary: $\int_{\Omega_{\bar{A}}} d\nu = \int_{\Omega_G} d\nu$.

[A20] A. Abrams. Extremal parameters and their duals for boundary maps associated to Fuchsian groups, *Illinois J. Math.* **65** (2021) No. 1, 153–179.

Lemma [B88, Appendix A3]: \forall oriented geodesic segment s on \mathbb{D} ,

$$\int_{\Psi^+(s)} d\nu = \text{length}(s),$$

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Since $\Omega_G = \bigcup_{i=1}^{8g-4} \mathcal{G}_i$, where

$$\mathcal{G}_k = \{(u, w) \mid uw \text{ exits } \mathcal{F} \text{ through side } k\} = \Psi_+(\text{side } k).$$

$$\int_{\Omega_G} d\nu = \sum_{k=1}^{8g-4} \int_{\mathcal{G}_k} d\nu = \sum_{k=1}^{8g-4} \text{length}(\text{side } k) = \text{Perimeter}(\mathcal{F}).$$

Thus $h_{\mu_A}(f_A) = \pi \cdot \frac{\text{Area}(\mathcal{F})}{\text{Perimeter}(\mathcal{F})}$.

[B88] F. Bonahon. The geometry of Teichmüller space via geodesic currents, *Inventiones Mathematicae* **92** (1988), 139–162.

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Fenchel–Nielsen coordinates on $\mathcal{T}(g) \approx \mathbb{R}^{6g-6}$ use a decomposition of S into $2g - 2$ pairs of pants by $3g - 3$ non-intersecting closed geodesics whose lengths can be manipulated independently.

One of them can be taken to correspond to one entire side of \mathcal{F} (blue) and can be made arbitrary long, i.e. $h_{\mu_{\bar{p}}}(f_{\bar{p}})$ can be made arbitrarily small.

Since perimeter of \mathcal{F} varies continuously within $\mathcal{T}(g)$, $h_{\mu_{\bar{A}}}(f_{\bar{A}})$ must take on all values between 0 and its maximum $H(g)$.

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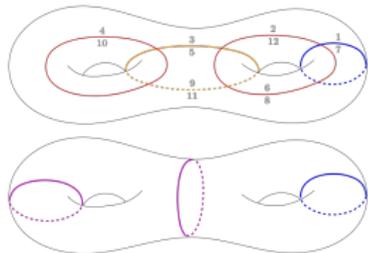
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Chain of $2g$ geodesics on S forming the sides of \mathcal{F} (top) and decomposition of S into $2g - 2$ pairs of pants by $3g - 3$ non-intersecting geodesics (bottom) for $g = 2$

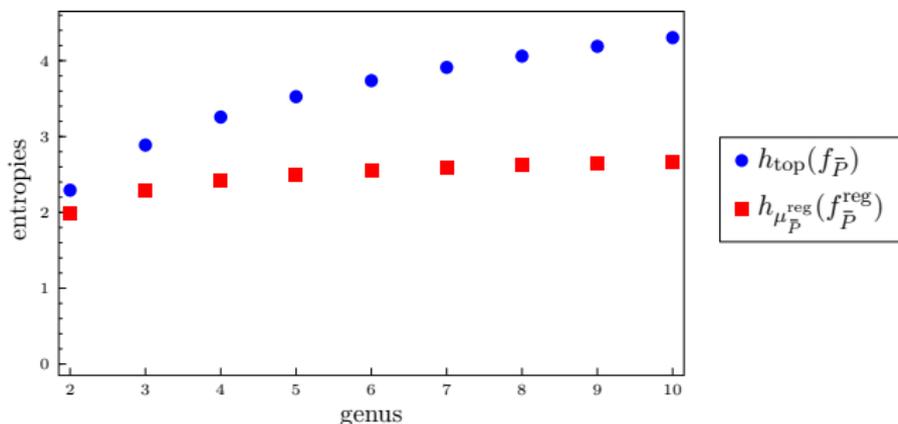
Comparison of entropies

Theorem: Smooth measure is not measure of maximal entropy.

If \bar{A} is extremal or has short cycles, the maximum of measure-theoretic entropy of $f_{\bar{A}}$ with respect to its smooth invariant measure $\mu_{\bar{A}}$,

$h_{\mu_{\bar{A}}^{\text{reg}}}(f_{\bar{A}}^{\text{reg}}) = \frac{\pi^2(4g-4)}{(8g-4) \operatorname{arccosh}(1+2 \cos \frac{\pi}{4g-2})}$ is strictly less than the topological

entropy of $f_{\bar{A}} = h_{\text{top}}(f_{\bar{A}}) = \log(4g - 3 + \sqrt{(4g - 3)^2 - 1})$ for all $g \geq 2$.



Topological entropy and maximum measure-theoretic entropy for different genera

Conjectures and question

- 1 For any $\mathcal{F} \in \mathcal{T}(g)$ and any \bar{A} with $A_k \in [P_k, Q_k]$, there exists a smooth $f_{\bar{A}}$ -invariant ergodic probability measure $\mu_{\bar{A}}$.

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- 2 For any \bar{A} with $A_k \in [P_k, Q_k]$, there exists a set $\Omega_{\bar{A}} \subset \mathbb{S} \times \mathbb{S} \setminus \Delta$ with finite rectangular structure that is a domain of bijectivity for $F_{\bar{A}}$ and moreover the global attractor of $F_{\bar{A}}$.

This is a part of the “Reduction Theory” for Fuchsian groups proposed by Don Zagier. Understanding the structure of $\Omega_{\bar{A}}$ may help in answering the following question:

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This is a part of the “Reduction Theory” for Fuchsian groups proposed by Don Zagier. Understanding the structure of $\Omega_{\bar{A}}$ may help in answering the following question:

- 3 Is it true that for any \bar{A} with $A_k \in [P_k, Q_k]$, the map $F_{\bar{A}}|_{\Omega_{\bar{A}}}$ is conjugate to F_{geo} by a map $\Phi_{\bar{A}} : \Omega_{\text{geo}} \rightarrow \Omega_{\bar{A}}$ that acts piecewise by Möbius transformations?

